

# CP violation in $\tau$ decay

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Belle Collaboration



# Outline

- 1 CP violation in SM
  - Understanding of CP asymmetry
  - CP Mechanism
- 2 CP violation in  $\tau$  decay
  - Ideas of CPV in  $\tau$
  - $\tau \rightarrow K \pi \nu$  decay
  - $\tau \rightarrow K \pi \pi \nu$  decay
- 3 status of  $\tau \rightarrow K_S h \pi^0$ 
  - Motivation
  - Event selection
  - Invariant Mass distribution



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# Origin of CP asymmetry

Electroweak and Yukawa coupling

$$\mathbf{V}_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

- Present known source of CP asymmetry is **CKM** matrix.
- CKM can be understood by Electroweak and Yukawa interaction of Higgs.
- Quark masses are not diagonalized in the flavor basis but diagonalized by unitary transformation. Then charged weak interaction are non-diagonal.
- 1 physical complex phase ( $\eta$ ) bring the CP violation in quark sector.



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# CP Mechanism

## General CP observable

CP violation arises from the inference

$$A_{X \rightarrow Y} = e^{i\theta_1} |A_1| e^{i\delta_1} + e^{i\theta_2} |A_2| e^{i\delta_2}$$

$$A_{\bar{X} \rightarrow \bar{Y}} = e^{-i\theta_1} |A_1| e^{i\delta_1} + e^{-i\theta_2} |A_2| e^{i\delta_2}$$

$\theta_i$  : CP phase,  $\delta_i$  : scattering phase (strong phase)

$$A_{CP} = \frac{|A_{X \rightarrow Y}|^2 - |A_{\bar{X} \rightarrow \bar{Y}}|^2}{|A_{X \rightarrow Y}|^2 + |A_{\bar{X} \rightarrow \bar{Y}}|^2} = \frac{-2|A_1||A_2| \sin(\theta_1 - \theta_2) \sin(\delta_1 - \delta_2)}{|A_1|^2 + |A_2|^2}$$

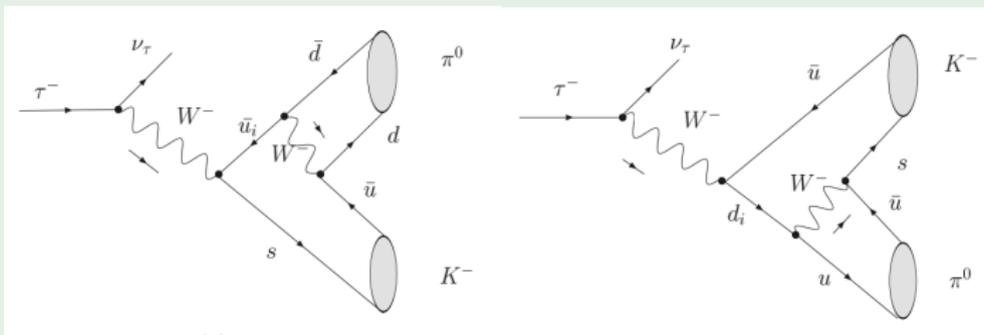


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## SM prediction of CP asymmetry in $\tau$ decays is too small to detect



Higher order contribution can only induce weak phase contribution (CKM)

$$A_{cp} = \frac{\Gamma(\tau \rightarrow K^+ \pi^0 \nu) - \Gamma(\tau \rightarrow K^- \pi^0 \nu)}{\Gamma(\tau \rightarrow K^+ \pi^0 \nu) + \Gamma(\tau \rightarrow K^- \pi^0 \nu)} \approx 2.3 \times 10^{-12}$$

what if we see the CP violation in  $\tau$  decays ?

This indicates you found the new physics so called "beyond the SM"



## Possible theory of CPV in $\tau$ decays

- R handed Vector Boson : L-R mixing
- Charged Higgs Doublet
- SUSY

In charged higgs model, higgs coupling is proportional to the mass difference of quarks.  $\rightarrow \Delta S \neq 0$  decays is promising in  $\tau$  CPV.

## Experimental efforts

- Search for CP violation in  $\tau \rightarrow K \pi \nu$  decays has been published by CLEO (2002).
- Same tryout is going on by  $\tau$  group in BELLE



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# introduction

## Situation

- CLEO has obtained the branching fraction with  $13.3 \text{ fb}^{-1}$  data.
- CLEO defined CP observable and set the limit on it with simple guess on structure function.
- BELLE has finished the study on structure function with rather complicated.



# Theory

Effective hamiltonian can be written by sum of SM and NP(new physics) term.

- $H_{eff} = H_{SM} + H_{NP}$
- $H_{SM} = \sin \theta_c \frac{G}{\sqrt{2}} \bar{\nu}_\tau \gamma_\mu (1 - \gamma_5) \tau \bar{s} \gamma^\mu (1 - \gamma_5) u + h.c$
- $H_{NP} = \sin \theta_c \frac{G}{\sqrt{2}} \eta_S \bar{\nu}_\tau (1 + \gamma_5) \tau \bar{s} u + h.c$

in the case of  $\tau \rightarrow K \pi \nu$ , hadronic current decays into two pseudo-scalar meson, which can be expanded in terms of the independent momenta  $(q_1 - q_2)^\mu$  and  $Q^\mu = q_1^\mu + q_2^\mu$

$$\bullet J_\mu = \langle h_1(q_1) h_2(q_2) | \bar{u} \gamma_\mu s | 0 \rangle = (q_1 - q_2)^\nu T_{\nu\mu} F(Q^2) + Q_\mu F_S(Q^2)$$

CP violating contribution incorporates to s-wave form factors by substitution:

$$\bullet F_S(Q^2) \rightarrow \tilde{F}_S(Q^2) = F_S(Q^2) + \frac{\eta_S}{m_\tau} F_H(Q^2)$$



# Theory

So, the amplitude for this channel,

$$M = \sin \theta_C \frac{G}{\sqrt{2}} \bar{u} \gamma_\mu (1 - \gamma_5) u [(q_1 - q_2)_\mu T^{\nu\mu} F + Q^\mu \tilde{F}_S]$$

The differential decay rate in hadronic rest frame,

$$\Gamma(\tau \rightarrow K \pi \nu_\tau) = M^2 d\Pi = \sum_X (L_X W_X) d\Pi$$

Where  $L_X$  and  $W_X$ , (for unpolarized  $\tau$ )

	$W_X$	$\bar{L}_X$
B	$4(q_1)^2  F ^2$	$\frac{1}{3}(2 + \frac{m^2}{Q^2}) - \frac{1}{6}(1 - \frac{m_{\tau}^2}{Q^2})(3 \cos^2 \psi - 1)(3 \cos^2 \beta - 1)$
SA	$Q^2  F_S ^2$	$\frac{m_\tau^2}{Q^2}$
SF	$4\sqrt{Q^2}  q_1  \text{Im}(F \tilde{F}_S^*)$	$-\frac{m_\tau}{Q^2} \cos \psi \cos \beta$
SG	$-4\sqrt{Q^2}  q_1  \text{Re}(F \tilde{F}_S^*)$	0



# Observable

## CP operation

$$d\Gamma_{\tau^-}(p_i, P, \eta_S) \rightarrow d\Gamma_{\tau^+}(-p_i, -P, \eta_S^*)$$

in order to get CP observable, subtract two the differential decay rates

## CP observable

$$\Delta(p_i) = \frac{d\Gamma_{\tau^-}}{d\Pi}(p_i) - \frac{d\Gamma_{\tau^+}}{d\Pi}(-p_i)$$

$$\Sigma(p_i) = \frac{d\Gamma_{\tau^-}}{d\Pi}(p_i, \eta_S = 0) + \frac{d\Gamma_{\tau^+}}{d\Pi}(-p_i, \eta_S^* = 0)$$

## CP observable

$$\xi^-(p_i) = \frac{\Delta(p_i)}{\Sigma(p_i)} = \frac{\bar{L}_{SF} \Delta W_{SF}}{\sum \bar{L}_X W_X}$$

by the way, CLEO defines somewhat different observable not including CP odd term in denominator

$$\xi^-(p_i) = \frac{\Delta(p_i)}{\Sigma(p_i)} = \frac{\bar{L}_{SF} \Delta W_{SF}}{\bar{L}_B W_B + \bar{L}_{SA} W_{SA}}$$



# Observable

## Averaged CP observable

The CP observable,  $\xi$ , can be integrated within particular range of momentum and  $Q^2$ .

$$\langle \xi^\pm \rangle = \int_{\Delta Q^2} \xi^\pm \frac{d\Gamma^\pm}{dQ^2 d\Omega} d\Omega dQ^2$$

## Characteristic of CP observable

- dependent on momentum and hadron invariant mass
- sensitive to structure function ( model dependence )  
 → different from particular decay mode



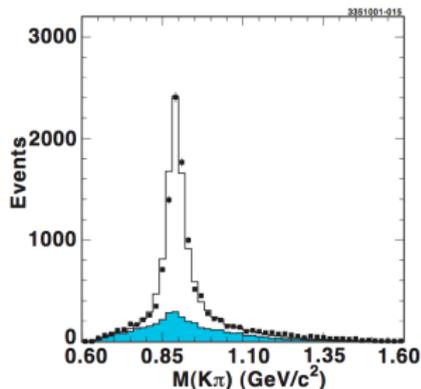
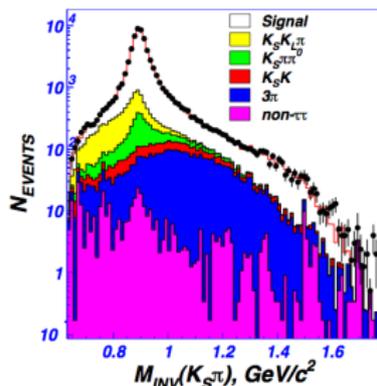
# Hadron structure in $\tau \rightarrow K_S \pi \nu$

## BELLE parametrization

- 2 solution for  $K_0^*(800) + K^*(892) + K_0^*(1430)$
- 1 solution for  $K_0^*(800) + K^*(892) + K^*(1410)$

## CLEO parametrization

CLEO used  $K^*(892) + K_0^*(1430)$



## MC generation

TAUOLA is MC generator for  $\tau$  decays which supports about 50 type of decays. For the decay  $\tau \rightarrow K \pi \nu$ ,  $F_S(Q^2)$  has been neglected  
 In order to study CPV with MC data, need another event generator  
 → Do we need to fix TAUOLA?

### event generation

It is possible to weight the event for some momentum and mass range by:

$$w = \frac{\sum(L_X W_X^{cp})}{\sum(L_X W_X^{tauola})}$$

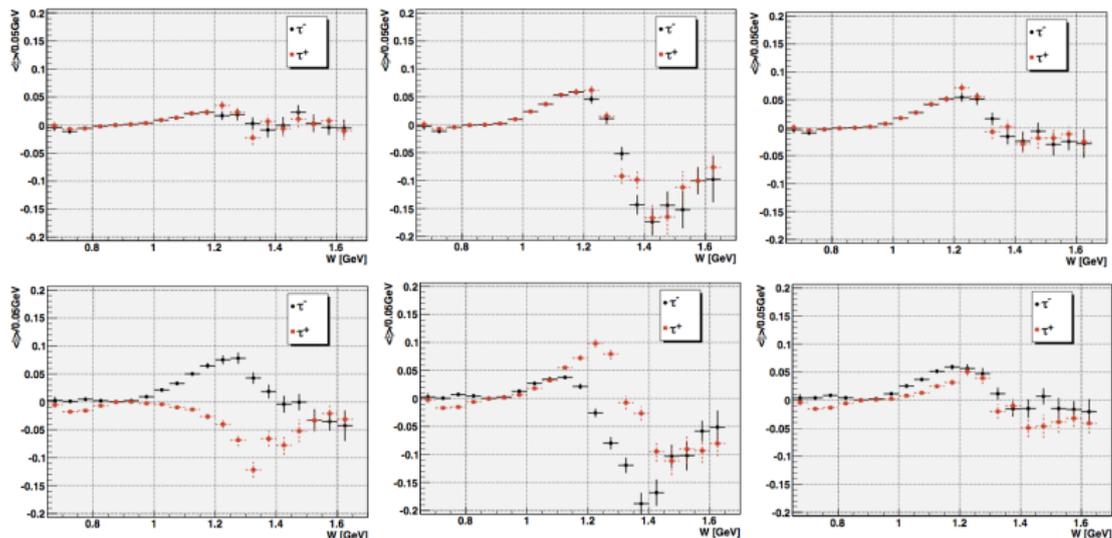
$W_X^{cp}$  is a function of  $\eta_S, s_1^2, s_2^2$  and  $Q^2$  so that don't need new generator.

This study has been done by Dr. Bischofberger and Prof. Hayashii (BELLE).



# MC study

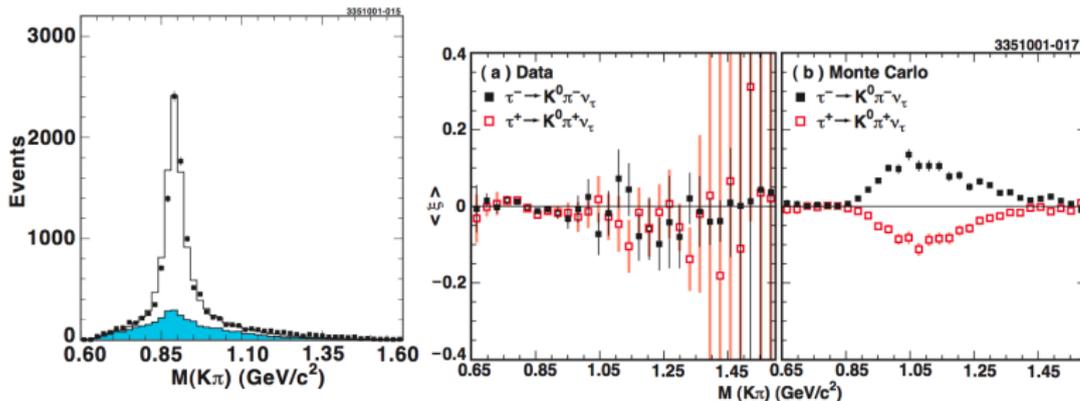
MC study for  $\langle \xi \rangle$  is done by using Belle parametrization result :



These result are obtained by assuming  $Im(\eta_S) = 1$ , maximal CP



# CLEO result and comparison



- CLEO limits (90 % C.L.)  $-0.172 < \Lambda < 0.067$  for  $13.3 \text{ fb}^{-1}$   
 $\Lambda$  is equivalent to  $\eta_S$  but use different normalization
- $F_S(Q^2) = 0$  and  $F_H = BW(K_0^*(1430))$



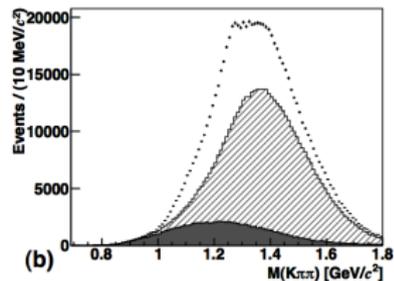
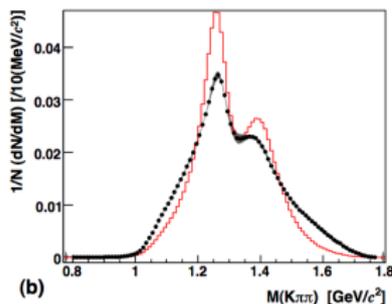
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# introduction

- No one has studied for 3 hadron case, because
  - No experimental information for the hadron structure for 3 hadron decay in  $\tau$  sector
  - seems rather complicated to obtain observable
- Dr. Lee Myung Jae has finished the study of this channel. but background contributes too much. cause problem of systematics in sensitivity.
- with successful extraction by unfolding method, we can see  $K_1(1270)$  and  $K_1(1400)$  clearly



# Theory

like as in  $\tau \rightarrow K \pi \nu$  decay,  $H_{NP}$  added to  $H_{SM}$  but,

$$H_{NP} = \sin \theta_c \frac{G}{\sqrt{2}} \eta_S \bar{\nu}_\tau (1 + \gamma_5) \tau \bar{s} u + \eta_P \bar{\nu}_\tau (1 + \gamma_5) \tau \bar{s} \gamma_5 u + h.c$$

$\eta_S$  is too small to contribute CP but  $\eta_P$  can contribute  
 Hadronic current is also different

$$\begin{aligned} J^\mu &= \langle K(p_1) \pi(p_2) \pi(p_3) | \bar{s} \gamma^\mu (1 - \gamma^5) u | 0 \rangle \\ &= [F_1(p_1 - p_3)_\nu + F_2(p_2 - p_3)_\nu] T^{\mu\nu} + i F_3 \epsilon^{\mu\nu\rho\sigma} p_{1\nu} p_{2\rho} p_{3\sigma} + F_4 Q^\mu \end{aligned}$$

- $F_1$  and  $F_2$  is due to  $K_1(1270)$  and  $K_1(1400)$  respectively.
- $F_3$  is the anomalous Wess-Zumino term.
- $F_4$  is scalar term, generally assumed to be negligible for this decay.



## Theory

Similarly to the previous work,

$$F_4 \rightarrow \bar{F}_4 = F_4 + \frac{F_H}{m_\tau} \eta_P$$

where pseudo-scalar form factor,  $F_H$ , has been defined :

$$F_H = \langle K(p_1) \pi(p_2) \pi(p_3) | \bar{s} \gamma^5 u | 0 \rangle$$



## observables in $\tau \rightarrow K \pi \pi \nu$ decay

3 kind of observable for 3 hadron decay. (general)

- Rate asymmetry

The difference between the rate for the process and that for the associated anti-process

In principle, same observable as for  $\tau \rightarrow K \pi$ .

- Polarization-dependent Rate asymmetry

Integrate decay rate for particular kinematic angle by weighting

- Triple-product Rate Asymmetry

each asymmetry observable is proportional to  $|F_H| \text{Im}(\eta_P)$ .



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## Motivation

- Branching fraction measurement (latest result from ALEPH,1998)
- Hadron Structure Function of 3 hadron in  $\tau$  decays.
- Study of Wess-Zumino anomaly term.
- CP violation ( for  $K_S \pi \pi^0 \nu$  decay )



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## Event selection

- Tauskim
- $2 \leq N_{trk} \leq 4$ , (at least 2 track from primary vertex)
- $20^\circ \leq \mathcal{M}_{missing} \leq 175^\circ$
- Thrust  $> 0.9$
- $E_{ECL} < 9.0$  GeV
- Event topology : 1-3
- Charge Sum = 0
- Particle ID LR cut :  $eid > 0.8$ ,  $muid > 0.8$ ,  $k/\pi > 0.7$
- $K_S^0$  :  $zdist < 2\text{cm}$ ,  $1.5\text{cm} < \text{flight length} < 40\text{cm}$
- $\pi^0$  :  $-6 < S_{gg} < 5$
- $E_\gamma < 0.3$  GeV

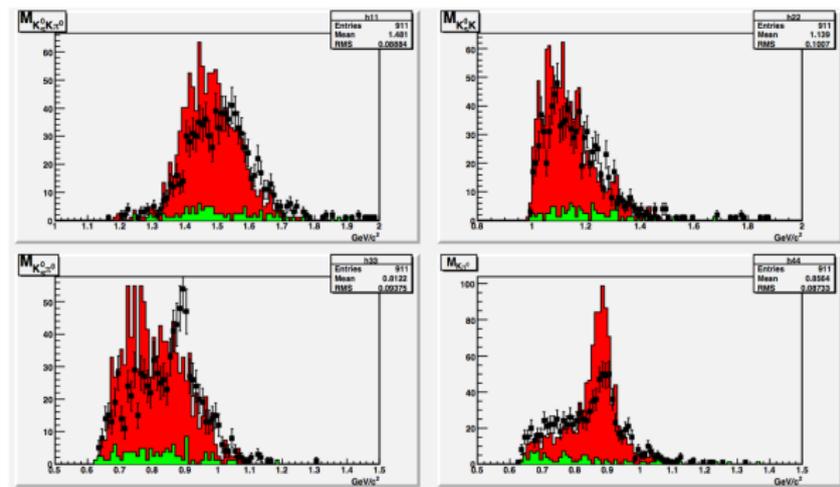


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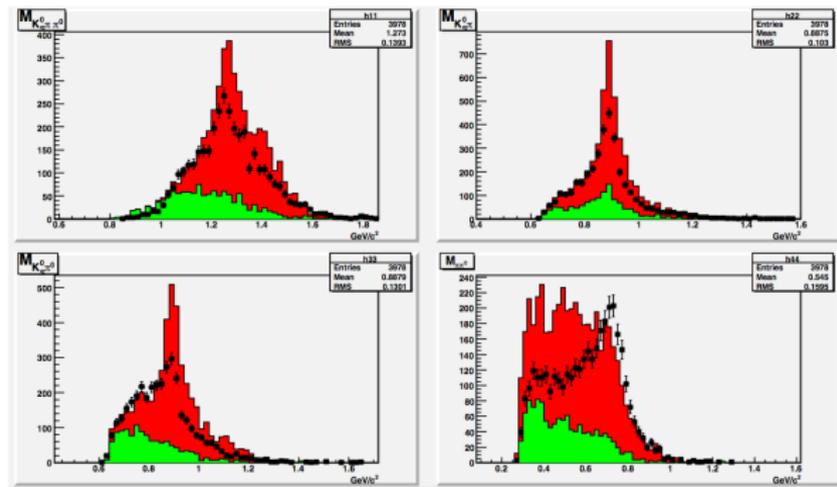
# Mass distribution for $\tau \rightarrow K_S K \pi^0 \nu$



- Statistics :  $67.7 \text{ fb}^{-1}$
- Vector current (  $K^*(892)$  and  $\rho(770)$  ) is seen
- expected background : 16.2 %  
 $\tau \rightarrow K_S K$  : 5.6 %,  $\tau \rightarrow K_S \pi^0$  : 7.6%, other  $\tau$  decays : 3.0%



# Mass distribution for $\tau \rightarrow K_S \pi \pi^0 \nu$



- Statistics : 67.7 fb<sup>-1</sup>
- Vector current (  $K^*$  (892) and  $\rho$ (770) ) is seen
- expected background : 34.6%

$\tau \rightarrow K_S K \pi^0$  : 6.2 %,  $\tau \rightarrow K_S \pi$  : 7.4 %,  $\tau \rightarrow K_S K_S \pi$  : 5.8 %,  $\tau \rightarrow \pi \pi^0$  : 5.8 %, other  $\tau$  decays : 10.0 %



## Summary

- CP violation in  $\tau$  decay is a definite sign of **New Physics**
- CLEO has published limits for  $\tau \rightarrow K^0 \pi \nu$  and  $\tau \rightarrow \pi \pi^0 \nu$  from an analysis of data corresponding to  $13.3 \text{fb}^{-1}$  and set the limit of CP violation by optimized CP observable  $\langle \xi \rangle$
- $\tau$  group in Belle has been studying CP violation with the same method carried out by CLEO.
- Not only for  $\tau \rightarrow K \pi \nu$  but also for  $\tau \rightarrow K \pi \pi \nu$  decay, CP observable can be studied by similar method.
- Before studying CP violation in  $\tau$  decay, one should understand the structure of the decay.
- Too much background affects  $\tau \rightarrow K \pi \pi \nu$  decay so that  $\tau \rightarrow K_S \pi \pi^0 \nu$  decay can be another option to be studied
- Study of  $\tau \rightarrow K_S \pi(K) \pi^0 \nu$  channel is on going.  
preliminary result of branching fraction will be obtained soon.



Thank you !



# For Further Reading I



A. Author.

*Physics in tau lepton.*

Some Press, 1990.



R. D. Kass

Search for CP violation in tau lepton decay

*Nuclear Physics B (Proc. Suppl.)*, 76 (1999) 215-218



Ken Kiers, etc.,

CP violation in  $\tau \rightarrow K\pi\pi\nu_\tau$

*arXiv:0808.1707v1*, [hep-ph] 12 Aug 2008



CLEO Collaboration

Search for CP violation in  $\tau \rightarrow K\pi\nu_\tau$  decays.

*Phys. Rev. Lett.* , 88 (2002) 111803-1



David Delepine

CP violation in Semi-Leptonic  $\tau$  decays

*arXiv:hep-ph/0702107v1*, 10 Feb 2007

