

EC induced by antineutrinos

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Introduction

- ν 's are one of the most abundant components of the universe.
- As well as the 2.7 K microwave background radiation, the universe is filled with a sea of relic ν 's that decoupled from the rest of the matter within the first 10s after the Big Bang.
- These relic ν 's may have played a crucial role in nucleosynthesis, structure formation and the evolution of the universe as a whole.

Introduction

- A non-zero neutrino mass can dramatically change the properties of the relic neutrino sea and its role in the evolution of the universe.
- The low E nuclear β or EC decays are excellent tools for testing non-zero ν masses.
- Among all β unstable nuclides, $^3\text{H}(\beta^-)$, $^{187}\text{Re}(\beta^-)$ & ^{163}Ho (EC) are attractive for (non-zero ν masses) due to their low Q-values.
- Recently the interest on the ^{163}Ho has been renewed.

Introduction

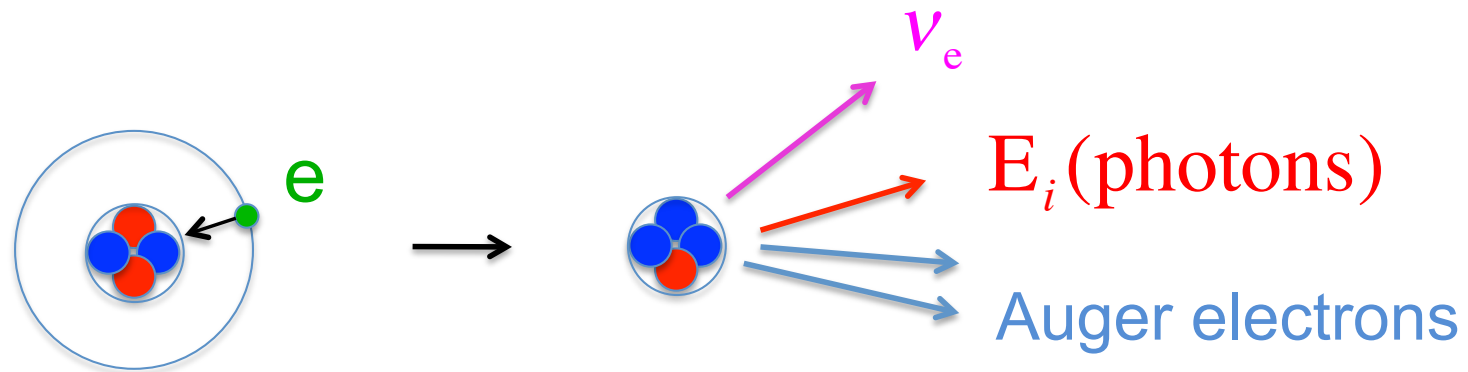
- M. Lusignoli & M. Vignati
“Relic anti- ν capture on ^{163}Ho decaying nuclei”
[PLB 697, 11 (2011) & PLB 701, 673 (2011) (E)]
- Y.F. Li & Z. Xing
“A possible detection of the cosmic anti- ν
background in the presence of flavor effects “
[PLB 698, 430 (2011)]

Introduction

- ‘Relic anti- ν captures on EC decaying nuclei’?
- Application to ^{163}Ho .
- Any other possible candidate nuclides?

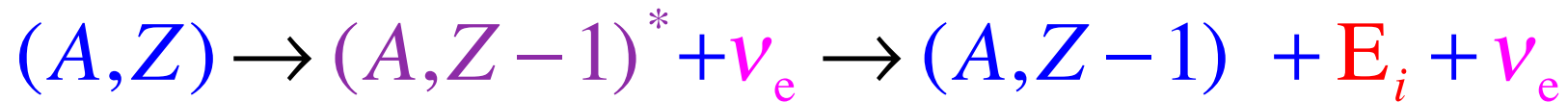
Electron Capture (EC) decay

EC : inverse beta decay



E_i : de-excitation energy of i -th atomic shell (K,L,M,..)
measurable!

Reaction rate for EC decay



$$\lambda_{EC} = \frac{G_\beta^2}{4\pi^2} \sum_i n_i C_i \beta_i^2 B_i (Q - E_i) \sqrt{(Q - E_i)^2 - m_\nu^2}$$

$$G_\beta = G_F \cos \theta_C$$

n_i : fraction of occupancy of the i - *th* atomic shell

C_i : nuclear shape factor

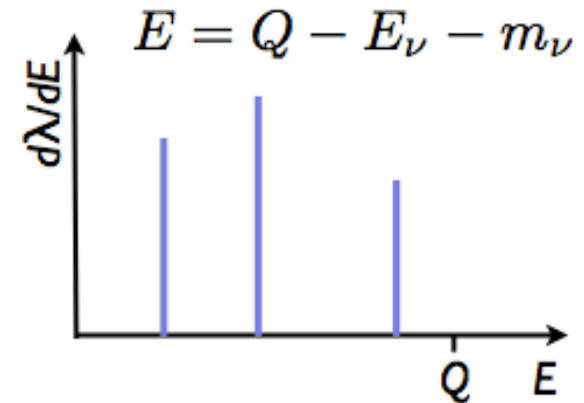
β_i : Coulomb amplitude of electron radial w.f.

B_i : atomic correction for electron exchange and overlap (~ 1)

Discrete energy spectrum

- In principle, the energy spectrum is given by a series of lines.

$$\frac{d\lambda_{EC}}{dE_c} = \frac{G_\beta^2}{4\pi^2} (Q - E_c) \sqrt{(Q - E_c)^2 - m_\nu^2} \times \sum_i n_i C_i \beta_i^2 B_i \delta(E - E_i)$$

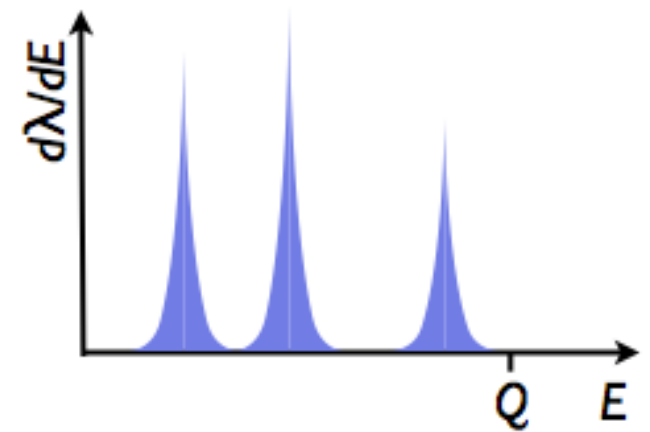


- The δ - ft. does not allow to measure the spectrum endpoint.
- EC is not sensitive to neutrino mass features at the endpoint.

Uncertainty relation

- Atomic levels have finite widths.
 → δ - ft. should be replaced by a **Breit-Wigner**.
 (De Rujula, Lusignoli, 1982)

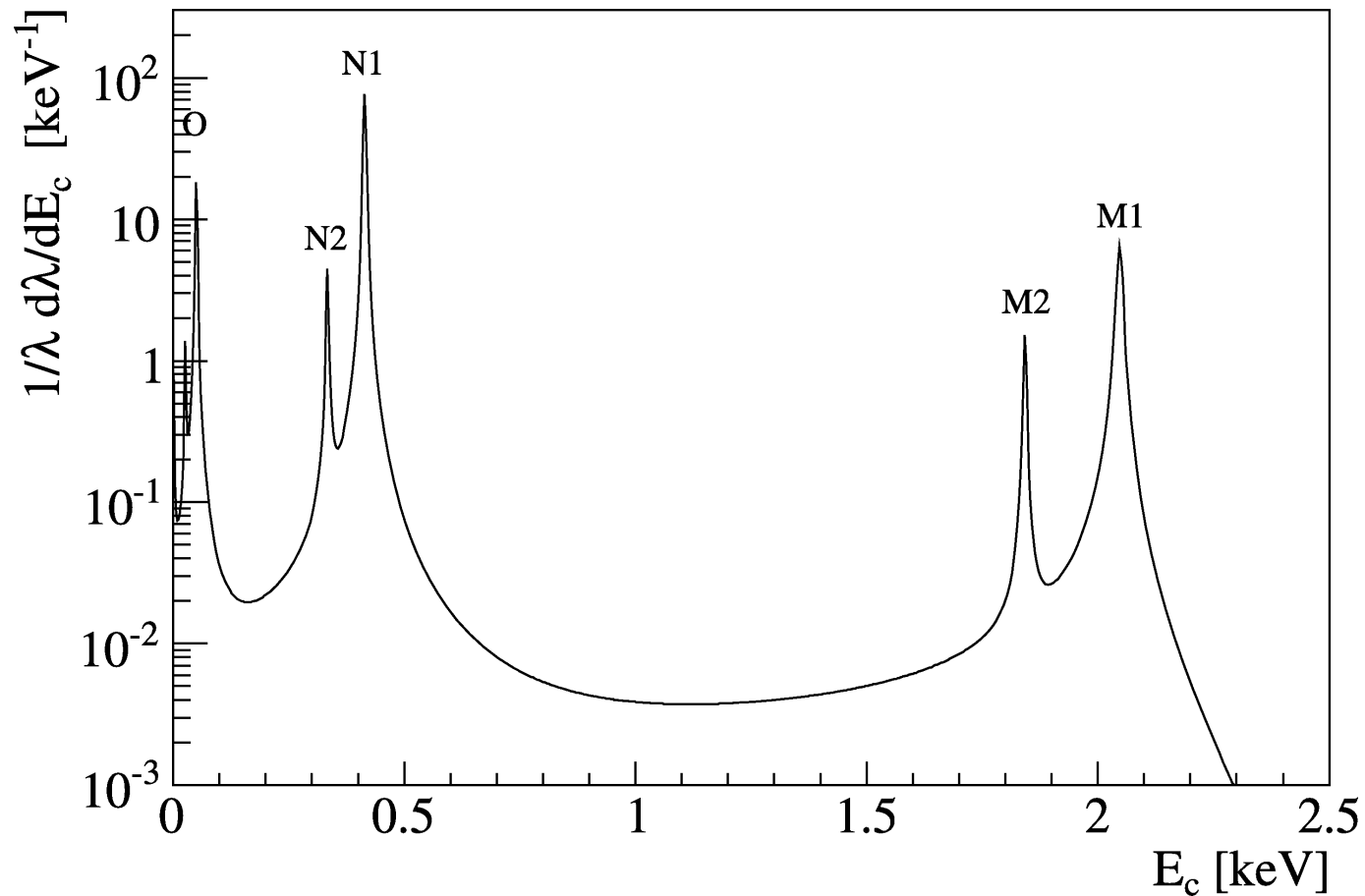
$$\frac{d\lambda_{EC}}{dE_c} = \frac{G_\beta^2}{4\pi^2} (Q - E_c) \sqrt{(Q - E_c)^2 - m_\nu^2} \times \sum_i n_i C_i \beta_i^2 B_i \frac{\Gamma_i}{2\pi (E_c - E_i)^2 + \Gamma_i^2 / 4}$$



- This allows E to reach the Q -value and to produce a neutrino with zero energy.
- Needs a nucleus with an E_i close to Q , to have small suppression from the Breit-Wigner.

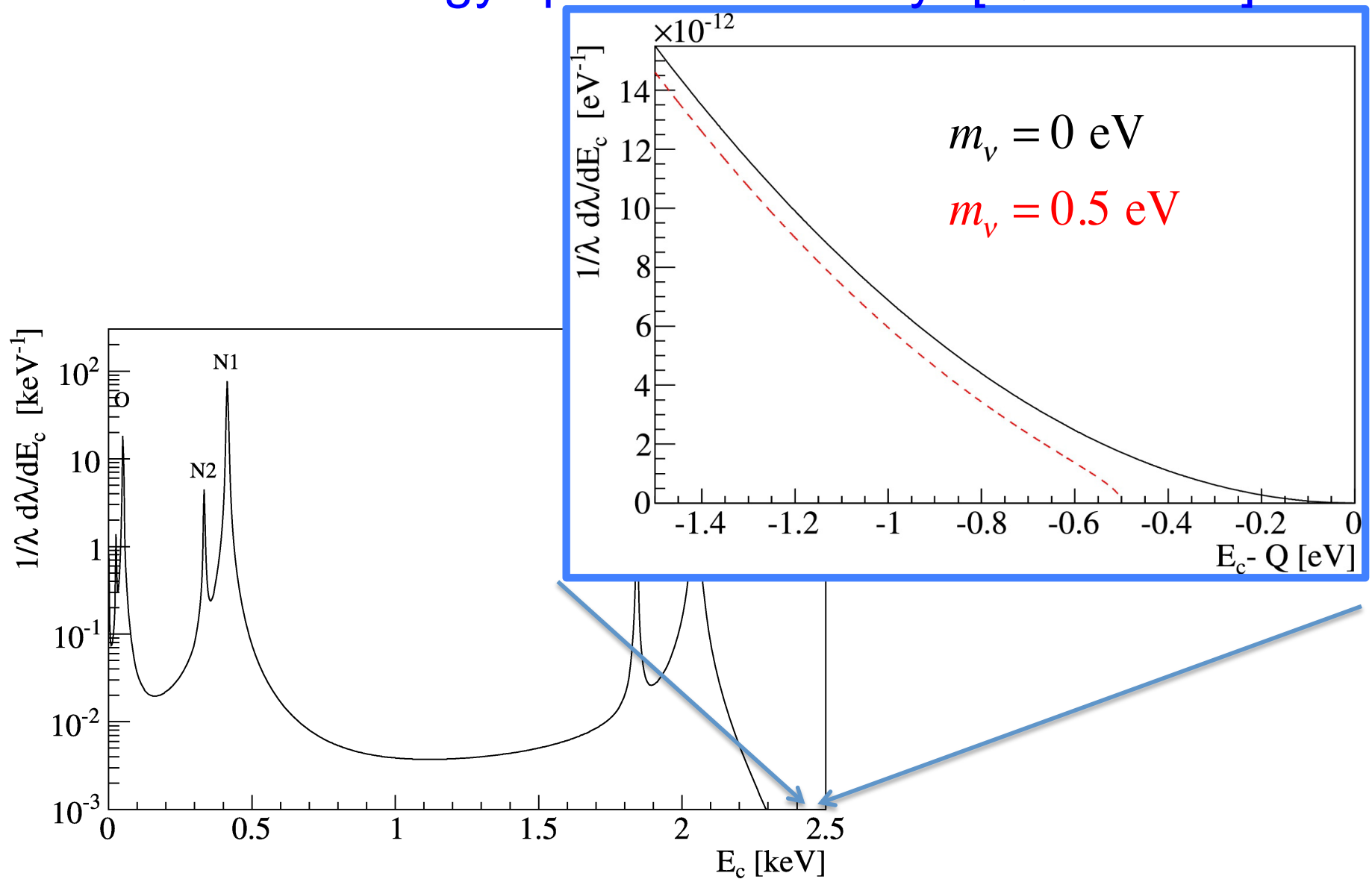
Energy spectrum in EC decays

De-excitation energy spectrum of $^{163}\text{Dy}^*$ [$Q=2.5$ keV]



Energy spectrum in EC decays

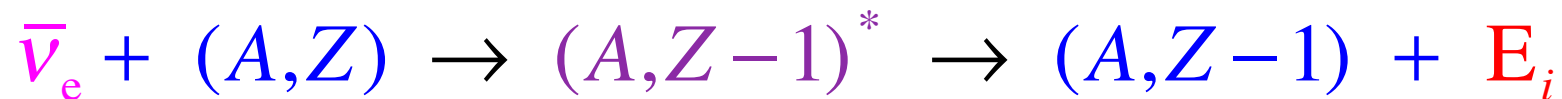
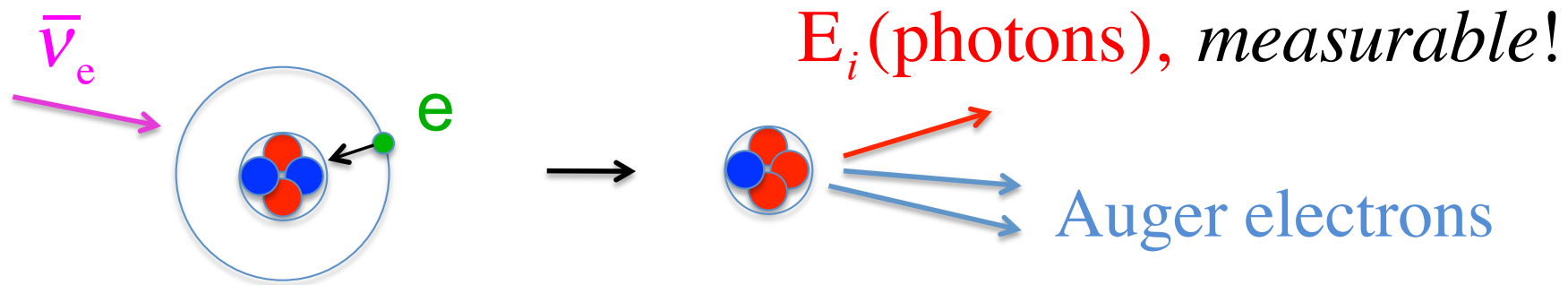
De-excitation energy spectrum of $^{163}\text{Dy}^*$ [$Q=2.5$ keV]



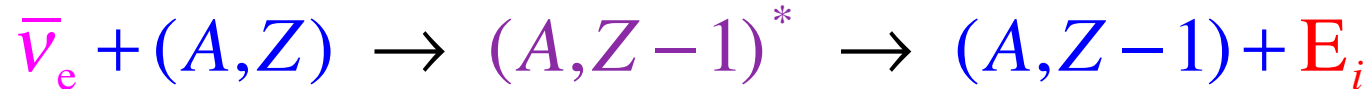
Relic $\bar{\nu}_e$ capture

Relic $\bar{\nu}_e$ capture on EC decaying nucleus

[No energy threshold $\rightarrow$ can absorb $\bar{\nu}_e$ with zero energy.]



Reaction rate for $\bar{\nu}_e$ capture



$$\lambda_{\bar{\nu}} = n_{\bar{\nu}} \frac{G_{\beta}^2}{2} \sum_i n_i C_i \beta_i^2 B_i \rho_i(E_{\bar{\nu}})$$

$$G_{\beta} = G_F \cos \theta_C$$

$n_{\bar{\nu}} (= n_{\nu})$: # density of incoming $\bar{\nu}$'s ($\sim 55/\text{cm}^3$)

$E_{\bar{\nu}}$: energy of incoming relic $\bar{\nu}$'s ($\simeq m_{\nu}$ for CνB)

β_i : Coulomb amplitude of electron radial w.f.

B_i : atomic correction for electron exchange and overlap (~ 1)

$\rho_i(E_{\bar{\nu}})$: density of final state per energy (Breit-Wigner form)

Interaction rate of relic $\bar{\nu}_e$ and EC

$$\frac{\lambda_{\bar{\nu}}}{\lambda_{EC}} \simeq 2\pi^2 n_{\bar{\nu}} \frac{\sum_i n_i \beta_i^2 B_i \rho_i(E_{\bar{\nu}})}{\sum_i n_i \beta_i^2 B_i (Q - E_i)^2}$$

$n_{\bar{\nu}} (= n_{\nu})$: # density of incoming $\bar{\nu}$'s ($\sim 55/\text{cm}^3$)

n_i : fraction of occupancy of i-th atomic shell

β_i : Coulomb amplitude of electron radial w.f.

B_i : atomic correction for electron exchange and overlap (~ 1)

$\rho_i(E_{\bar{\nu}})$: density of final state per energy

Interaction rate for ^{163}Ho

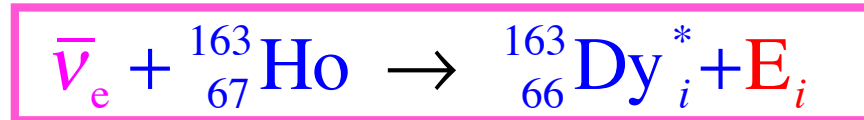
- ^{163}Ho

- $T_{1/2} : 4570 \text{ y}$
- Q-value for EC : 3.00 keV (NNDC at BNL)
[* Recent : $2.80 \pm 0.05 \text{ keV}$ (Gomez, 2008)]

- EC decay in ^{163}Ho



- Relic $\bar{\nu}_e$ capture on ^{163}Ho



Application to ^{163}Ho

Energy levels & their widths of the captured electrons, for ^{163}Dy .

Level	E_i (eV)	Γ_i (eV)
M_1 [3S]	2047	13.2
M_2 [3P $_{1/2}$]	1842	6.0
N_1 [4S]	414.2	5.4
N_2 [4P $_{1/2}$]	333.5	5.3
O_1 [5S]	49.9	3
O_2 [5P $_{1/2}$]	26.3	3
P_1 [6S]	0.0	1

Application to ^{163}Ho

Electrons squared wave functions at the origin β_i^2 relative to β_{3S}^2
[I.M. Band et al., At. Data. Nucl. Data Tables 35, 1 (1986).]

Levels	Ratio
$3P_{1/2} / 3S$	0.0526
$4S / 3S$	0.232
$4P_{1/2} / 3S$	0.0119
$5S / 3S$	0.0345
$5P_{1/2} / 3S$	0.0015
$6S / 3S$	0.0021

For given principle Q.#.,
electron squared w. f. of S-state ($l=0$) is ~ 20 times larger
than the one of $P_{1/2}$ state

Application to ^{163}Ho

Interaction rate for ^{163}Ho

$$\frac{\lambda_{\bar{\nu}}}{\lambda_{EC}} \simeq 2\pi^2 n_{\bar{\nu}} \frac{\sum_i n_i \beta_i^2 B_i \rho_i(E_{\bar{\nu}})}{\sum_i n_i \beta_i^2 B_i (Q - E_i)^2}$$

Q-value	2.2 keV	2.5 keV	2.8 keV
$\lambda_{\bar{\nu}} / \lambda_{EC}$	7.6×10^{-22}	5.8×10^{-23}	1.4×10^{-23}

Feasibility

- If there is no overdensity $\left(= \frac{\rho(x) - \bar{\rho}}{\bar{\rho}} \right)$ in the ν distribution, the detection is not possible.
- Kaboth et al., [PRD82, 2010, KATRIN experiment]
“If the overdensity is $> 2 \times 10^9$,
KATRIN experiment may be able to detect relic neutrinos”

Any other candidates?

Information for searching

- 'stable.dat' : A, element, Z, mass excess, J^π , $T_{1/2}$, decay mode, decay branch, decay Q-value.
- 'nudat78.tab' : abundance
- 'element-abundance.dat' : element abundance
- 'e_BE.dat' : electron binding energy

Any other candidates?

Conditions for candidates for resonant EC

- Q-value < 10 keV
- Abundance < 0.1% (should exist)
- $T_{1/2}$ of daughter nuclei (z-1) < 100 y
- Selection rule allowed : $\Delta J=1$, $\Pi_f \Pi_i = +1$

Any other better candidates?

Resonant Capture of orbital electron of K & L shells

G->G : L-shell EC decay

A	Z	life(z)	d-mode	jpia	abun(z)	Q_EC (keV)
157 TB	65	0.2239E+10	EC	3/2+	0.000000	60.00
163 HO	67	0.1441E+12	EC	7/2-	0.000000	3.00
187 OS	76	0.0000E+00		1/2-	1.600000	-3.00
193 PT	78	0.1577E+10	EC	1/2-	0.000000	57.00
194 HG	80	0.1400E+11	EC	0+	0.000000	40.00
202 PB	82	0.1656E+13	A	0+	0.000000	50.00
205 PB	82	0.4825E+15	EC	5/2-	0.000000	51.00
213 AT	85	0.1250E-06	A	9/2-	0.000000	73.00
215 RN	86	0.2300E-05	A	9/2+	0.000000	82.00
243 CM	96	0.9177E+09	EC	5/2+	0.000000	8.90

G->G : K-shell EC decay

157 TB	65	0.2239E+10	EC	3/2+	0.000000	9.76
235 NP	93	0.3422E+08	A	5/2+	0.000000	8.09
243 CM	96	0.9177E+09	EC	5/2+	0.000000	6.90

Any other candidates?

- Krauss & Grashow, Nature 310, 19(1984) & others
→ do not consider **isomer (meta-stable state)**
- Meta-stable
 - $T_{1/2} > 10^{-9} \text{ s}$
 - Occasionally, $T_{1/2}$ is far longer.
[lasting minutes, hours, or 10^{15} years for $^{180\text{m}}\text{Ta}$.]

Calculation of Q-value for EC with allowed selection rule
for **all electron states** and **meta-stable states**

Conclusions

- EC decaying nuclei can capture relic anti- ν 's.
- However, if there is **no overdensity** in the ν distribution, the detection is **not possible**.
- Kaboth et al., “If the overdensity is $> 2 \times 10^9$, KATRIN experiment may be able to detect relic ν 's.”
- **Any other possible candidate nuclides** with resonant EC decaying nuclides for CvB detection (considering **meta-stable states & all electron states**)