

**Review on**  
**Proposed antimatter gravity measurement with an antihydrogen beam**

A. Kellerbauer, Nuclear Instruments and Methods in Physics Research B 266(2008)

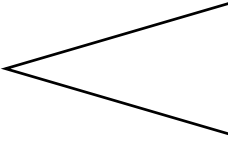
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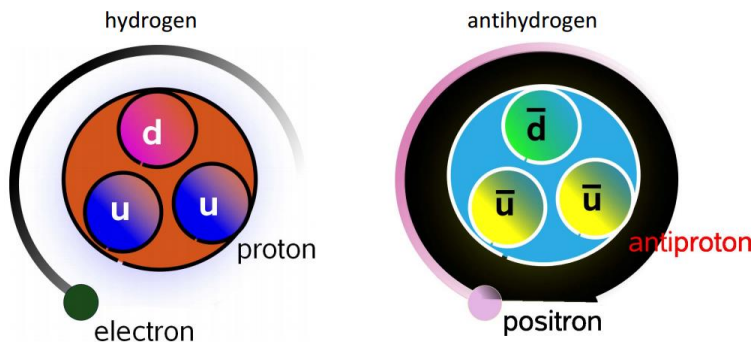
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1. Motivation
  2. AEGIS experiments
    - a. Aims
    - b. Methods
  3. Conclusion

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# 1. Motivation

# 1. Motivation

- Modern physics 
  - General relativity
    - about antimatter
    - Weak Equivalence Principle
  - Quantum mechanics
    - CPT symmetry  
(matter-antimatter symmetry)
- Antimatter : a material composed of anti-particles



1928, prediction of antimatter (Dirac)  
1932, discovery of positron in cosmic rays (Anderson)  
1955, antiproton discovery (Segre)  
1956, antineutron discovery (Cork)  
1995, creation of high-energy antihydrogen (CERN)  
2002, creation of 10K antihydrogen (ATHENA, ATRAP)  
2011, antihydrogen confinement (ALPHA)

Marco G., AEGIS at CERN : measuring Antihydrogen fall, INFN, Posmol(2011)

# 1. Motivation

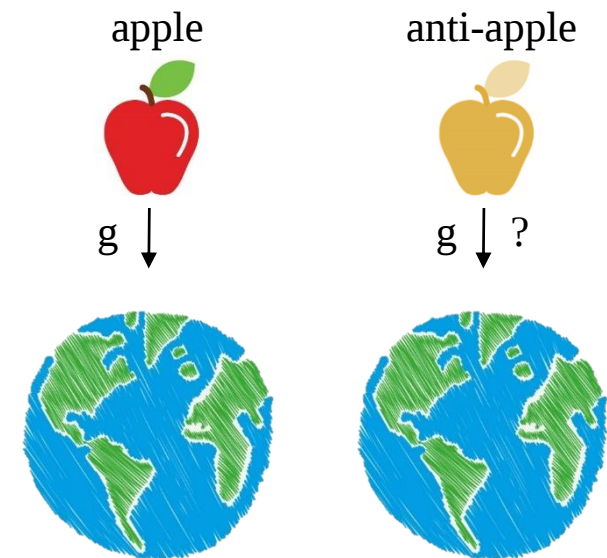
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**Weak equivalence principle (WEP)**, in the general relativity

“ In a uniform gravitational field,  
all objects fall with the same acceleration,  
regardless of their composition.”

WEP is tested with matter to high accuracy :  $\sim 10^{-16}$  (2016)

But the behavior of anti-matter in a gravitational field  
has never been tested experimentally.



$$m_g = \overline{m_g}$$

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# 1. Motivation

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## Weak equivalence principle(WEP), in the general relativity

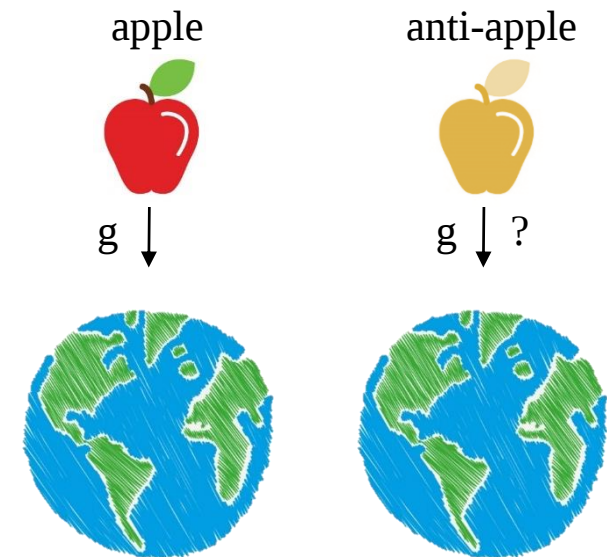
Gravity is **the only force** not described by a quantum field theory.

According to a hypothetical quantum theory of gravity,  
gravity could have with a tensor, vector, and scalar component.  
A vector component might lead a repulsive force.

That constitute a **violation of the WEP**.

Anti-gravity can be a possible solution for cosmological problem.  
But still a number of arguments against.

In order to repair contradictions and to incorporate gravity into the Standard model,  
observing the behavior of antimatter in gravitational field is necessary.



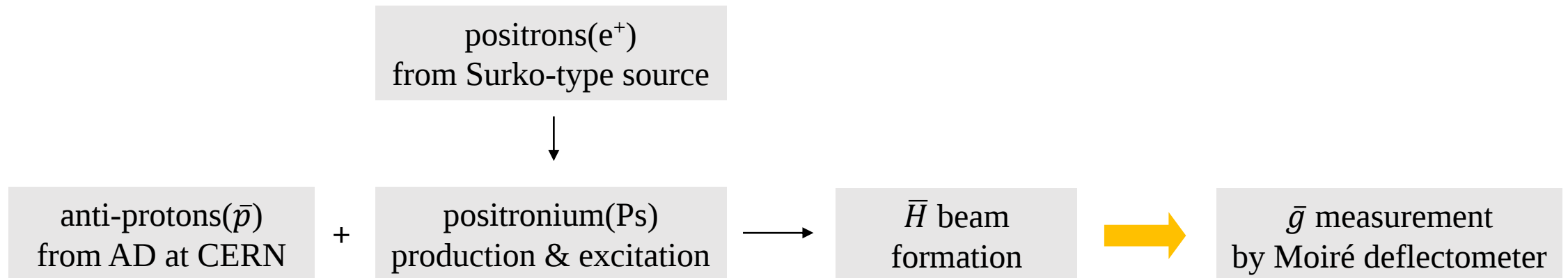
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## 2. AEGIS experiments – Aims

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**the AEGIS(Antimatter Experiment : Gravity, Interferometry, Spectroscopy) experiments**

: the direct measurement of the Earth's local gravitational acceleration  $\bar{g}$  on  $\bar{H}$  with 1% precision

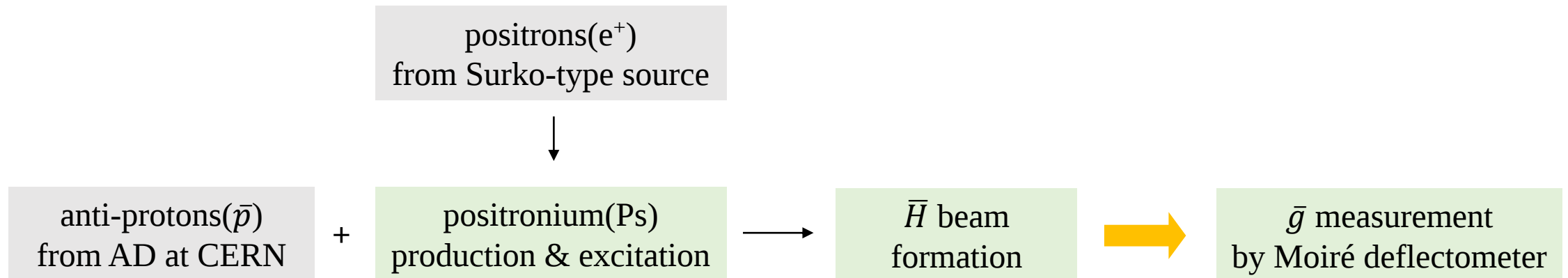


## 2. AEGIS experiments – Aims

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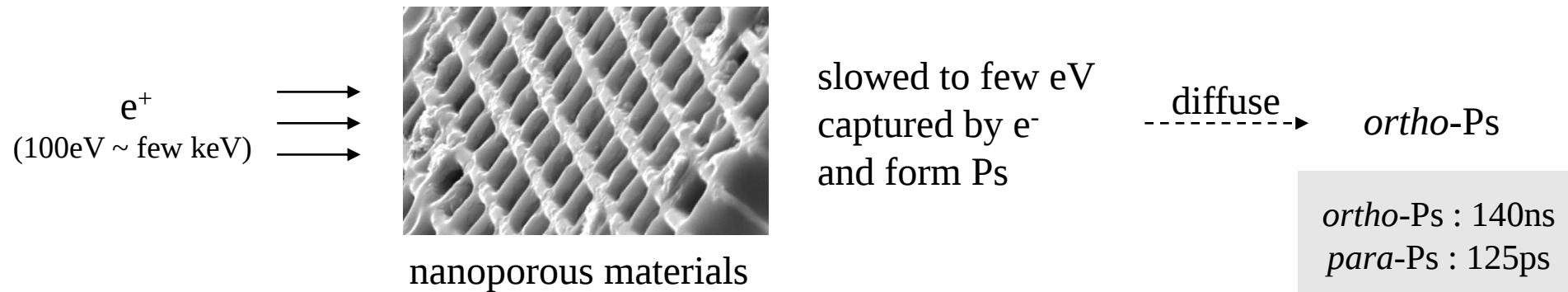


## 2. AEGIS experiments – Methods

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### ① Ps production and excitation

A sizable fraction of Ps diffuses out of the nanoporous material film.



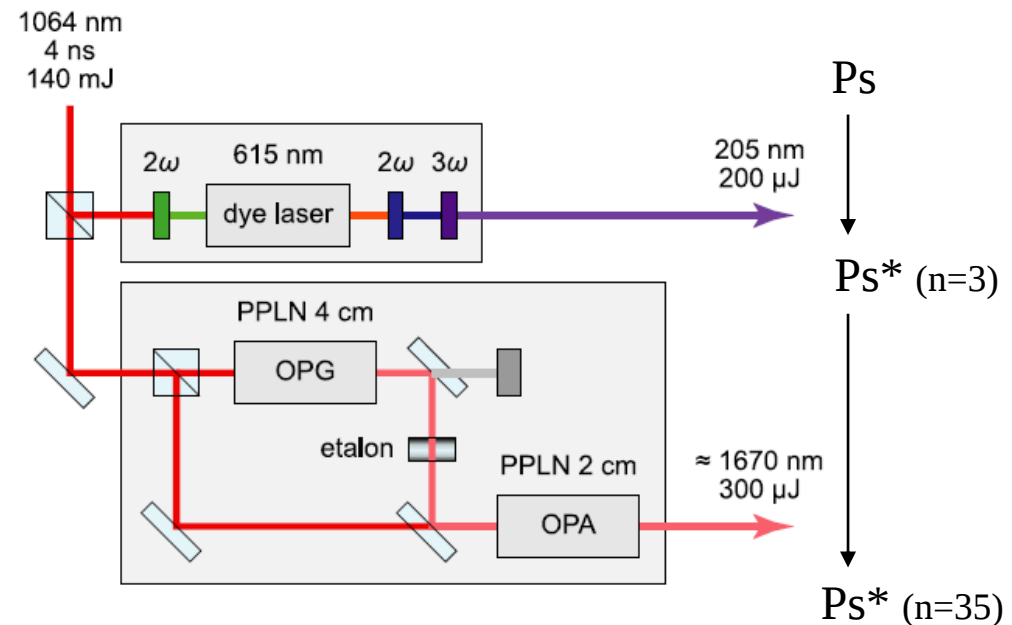
The yield depend on the target material, temperature and the implantation depth.  
(~30% in silicon-based polymer materials cooled to 50K)

## 2. AEGIS experiments – Methods

### ① Ps production and excitation

The cross-section of the  $\bar{H}$  formation reaction :  $\sigma \approx a_0 n^4 \rightarrow$  excitation of Ps is needed

The Ps excited to a Rydberg state with principal quantum number  $n=30\dots40$  by two pulsed-laser systems.

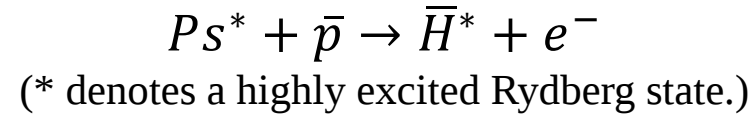


## 2. AEGIS experiments – Methods

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### ② $\bar{H}$ beam formation

The charge exchange reaction



- large cross-section ( $\approx 10^7 \text{\AA}^2$  for  $n = 35$ )
- narrow and well-defined band of  $\bar{H}^*$
- determination of velocity distribution of  $\bar{H}^*$  by  $\bar{p}$  temperature ( $\sim 25 - 80 \text{ m/s}$ )

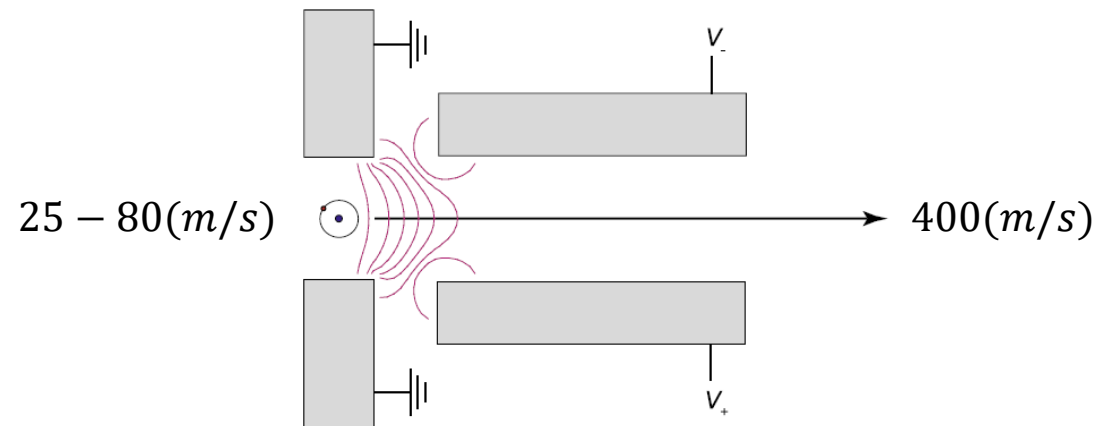
## 2. AEGIS experiments – Methods

### ② $\bar{H}$ beam formation

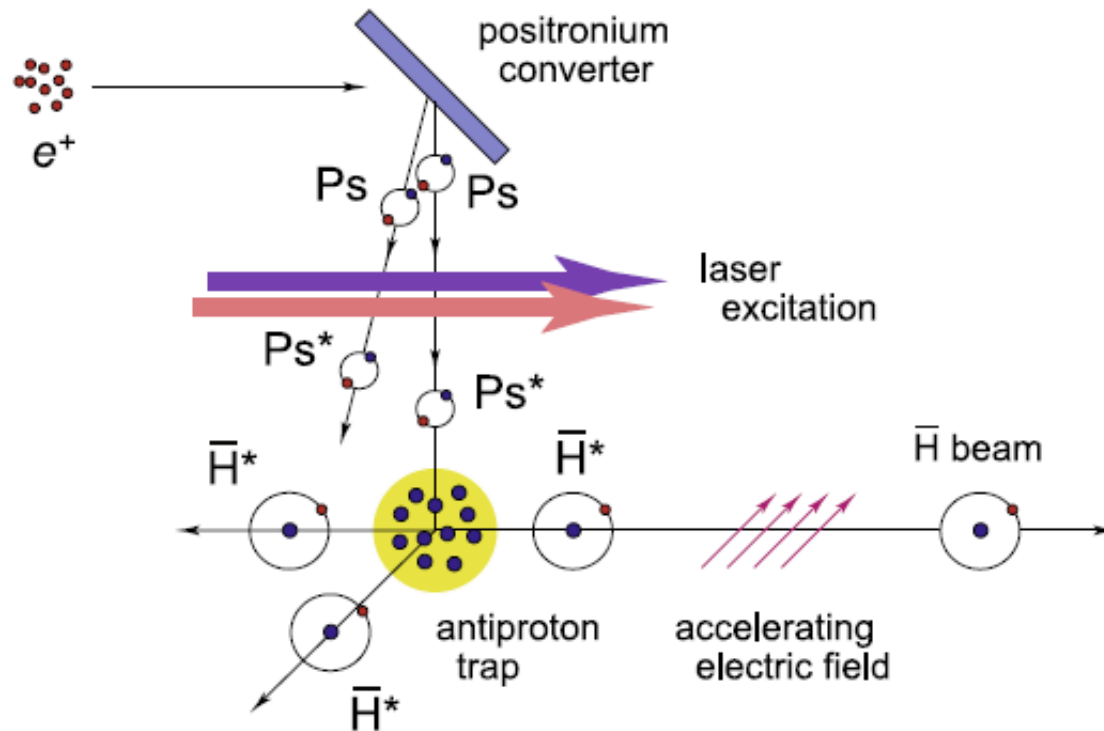
#### Stark acceleration

Atoms are accelerated or decelerated (experienced the force) by electric field gradient.  
Rydberg atoms which have large dipole moment are amenable.  
they are accelerated into a beam by inhomogeneous electric field.

$$\vec{F} = -\frac{2}{3}ea_0n(n-1)\nabla\vec{E}$$



## 2. AEGIS experiments – Methods



< Principle sketch until now >

1.  $e^+$  capture
2.  $\bar{p}$  production
3.  $Ps$  formation
4.  $Ps$  excitation
5.  $\bar{H}$  production
6.  $\bar{H}$  beam formation

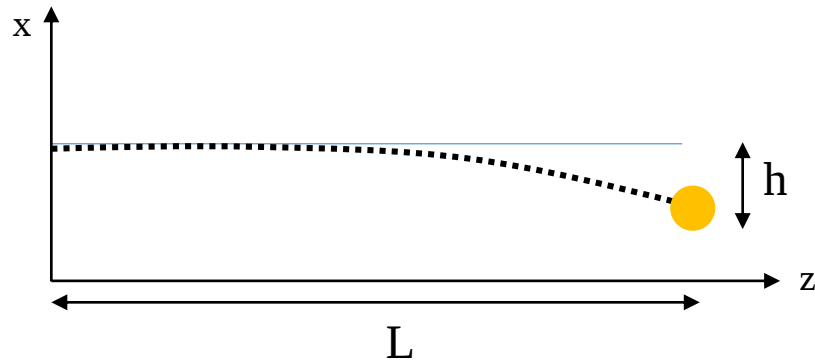
Note :  $\bar{p}$  trap before acceleration

## 2. AEGIS experiments – Methods

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### ③ Gravity measurement

Free-fall of  $\bar{H}$  cannot be detected in usual way



$$h = \frac{1}{2} g T^2 = \frac{1}{2} g \left( \frac{L}{v_z} \right)^2$$

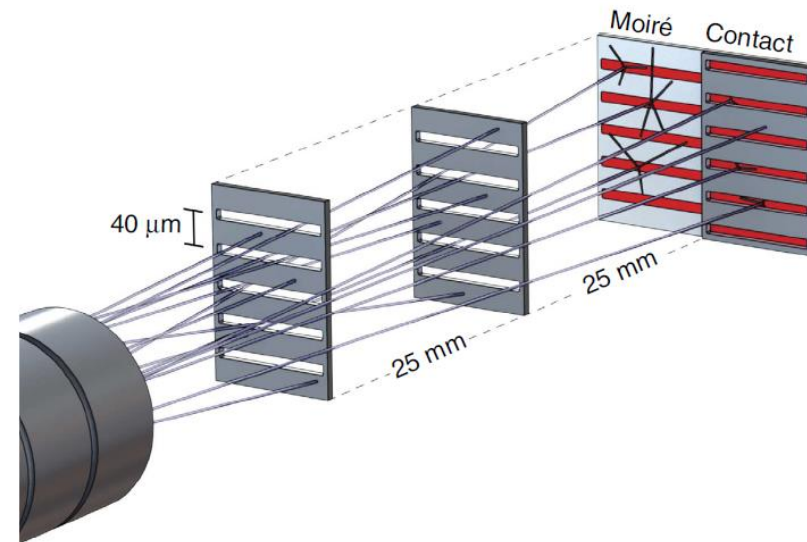
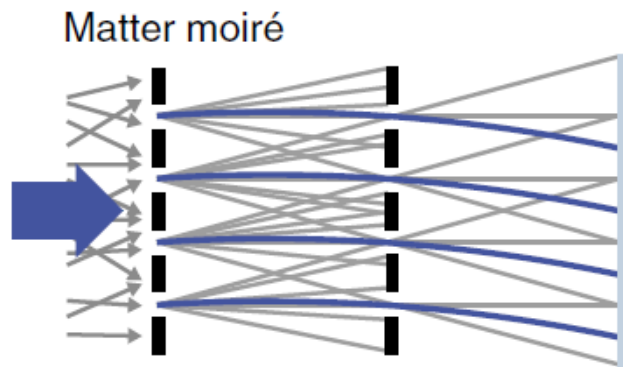
*T : time of flight*

$h \sim 20\mu\text{m}$ , where  $L \sim 1\text{m}$ ,  $v_x \sim 500\text{m/s}$   
much larger than beam radial size after 1m flight ( $\sim \text{cm}$ )

## 2. AEGIS experiments – Methods

### ③ Gravity measurement

Free-fall detection by Moiré deflectometer

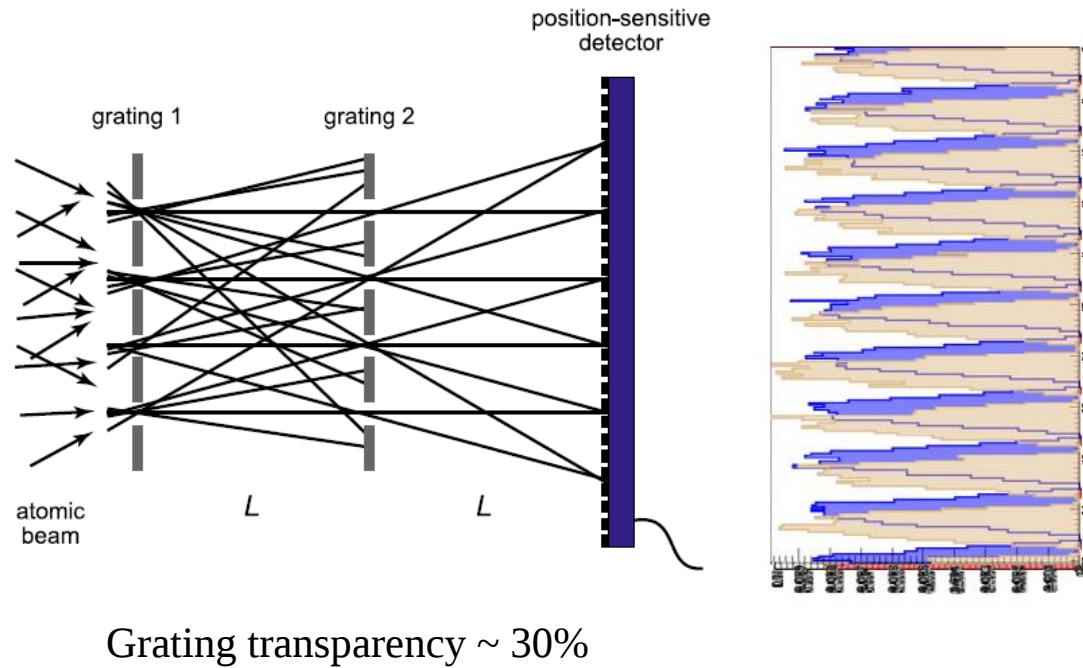


S. Aghion, A moiré deflectometer for antimatter, nature communications(2014)

## 2. AEGIS experiments – Methods

### ③ Gravity measurement

Free-fall detection by Moiré deflectometer



$a$  : the grating period ( $\sim 80\mu m$ )  
 $d$  : the gap in gratings  
 $L$  : the length between gratings ( $\sim 40cm$ )

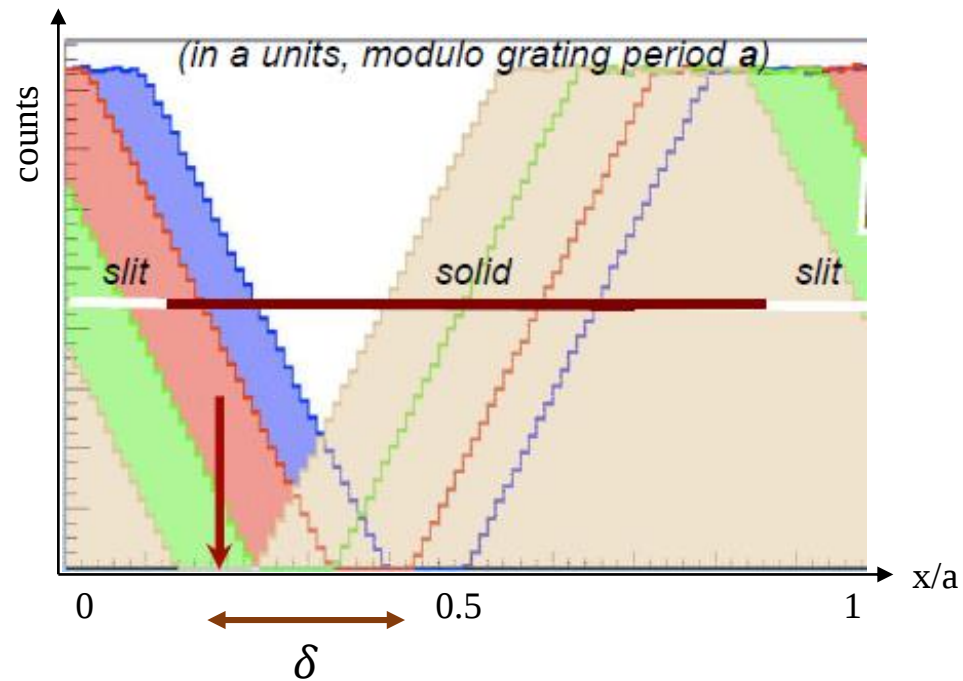
It is NOT a quantum deflectometer.  
No diffraction occurs.

$$\frac{d^2}{L} \gg \lambda_{dB}$$

## 2. AEGIS experiments – Methods

### ③ Gravity measurement

Free-fall detection by Moiré deflectometer



$$\delta = 2 * \left( \frac{1}{2} g T^2 \right) / a = \frac{g T^2}{a}$$

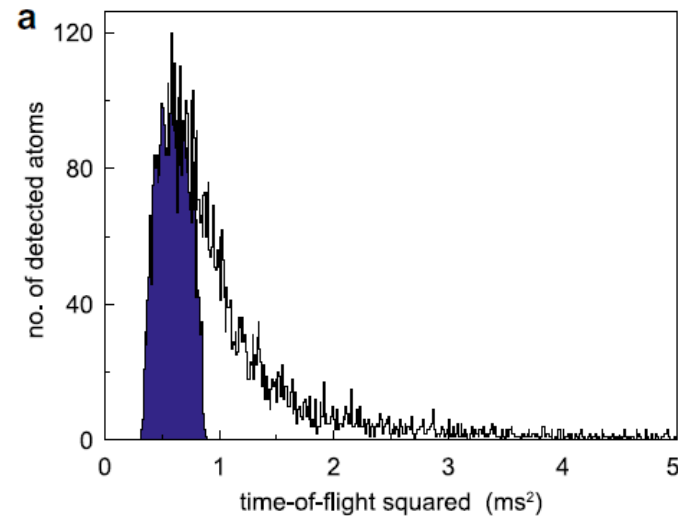
## 2. AEGIS experiments – Methods

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### ③ Gravity measurement

Extraction of  $g$  value from observables( $T, \delta$ ) – illustrated by Monte-carlo simulations

First, record vertical position for each event as a function of  $T$ (or  $v$ )



And bin these events in symmetric classes of  $T^2$

$$v_{beam} = 600 - 550 - 500 - \dots - 200 \text{ m/s}$$

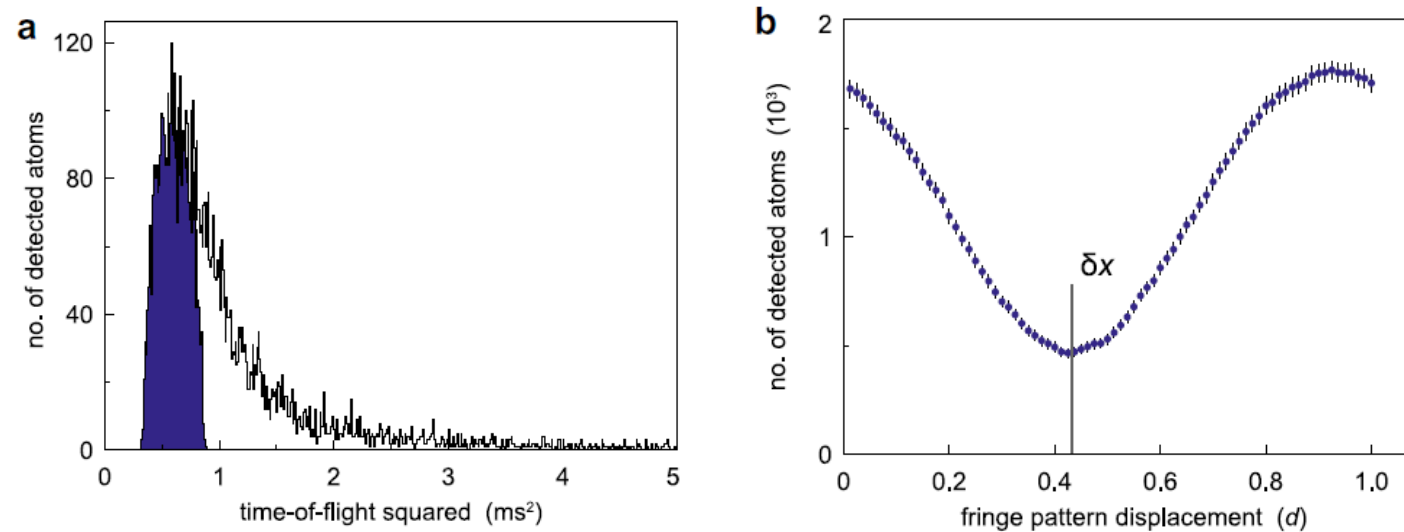
## 2. AEGIS experiments – Methods

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### ③ Gravity measurement

Extraction of  $g$  value from observables( $T, \delta$ ) – illustrated by Monte-carlo simulations

Next, extract  $\delta x$  of the fringe pattern for each of the count classes

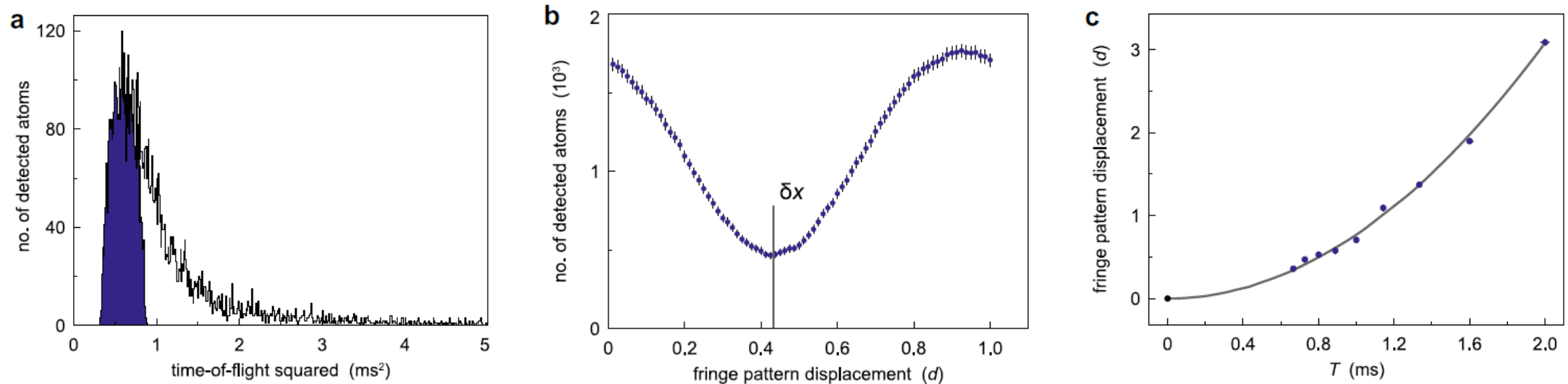


## 2. AEGIS experiments – Methods

### ③ Gravity measurement

Extraction of  $g$  value from observables( $T, \delta$ ) – illustrated by Monte-carlo simulations

Finally, plot  $\delta x$  against the mean time of flight  $\sqrt{\langle T^2 \rangle}$



## 2. AEGIS experiments – Methods

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### ③ Gravity measurement

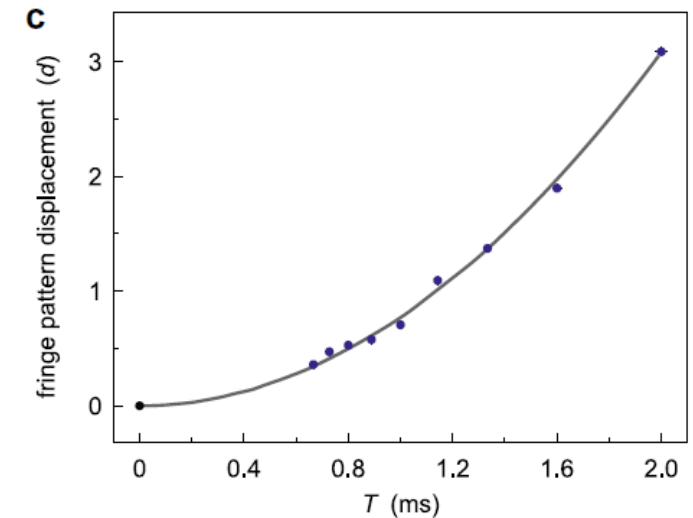
Extraction of  $g$  value from observables( $T, \delta$ ) – illustrated by Monte-carlo simulations

Finally, plot  $\delta x$  against the mean time of flight  $\sqrt{\langle T^2 \rangle}$

the coefficient of a quadratic fit to that yields  $\bar{g}$

$$\delta x = \delta * a = gT^2$$

(1% relative precision, about  $10^5 \bar{H}$  atoms at 100mK)



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- **AEGIS** =  $\begin{array}{c} \text{ATHENA} \\ \text{ATRAP} \\ \text{(at the AD)} \end{array}$  +  $\begin{array}{c} \text{Charge exchange production of } \bar{H} \\ \text{Stark acceleration} \\ \text{The Moiré deflectometer} \end{array}$
- Obtaining  $\bar{H}$  atoms at 100mK is an essential part, since the precision of gravity measurement is mainly limited by the  $\bar{H}$  temperature.
- If even colder  $\bar{H}$  ensembles can be attained, gravity measurements with even higher precision as well as competitive CPT tests through spectroscopy.